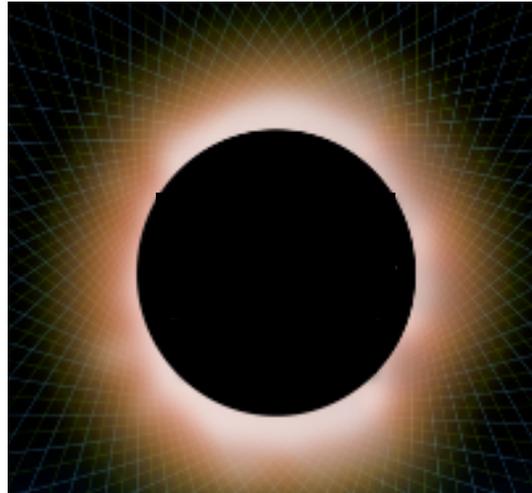


# Black Holes and Quantum Mechanics



Juan Maldacena

Institute for Advanced Study

Letter from Schwarzschild to Einstein.

22 December 1915,

$$ds^2 = \left(1 - \frac{\gamma}{R}\right) dt^2 - \frac{dR^2}{\left(1 - \frac{\gamma}{R}\right)} - R^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

It as been very confusing ever since...

Classically...

$$ds^2 = \underbrace{\left(1 - \frac{\gamma}{R}\right)} dt^2 - \frac{dR^2}{\left(1 - \frac{\gamma}{R}\right)} - R^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

Extreme slow down of time

# The coordinate ``singularity''

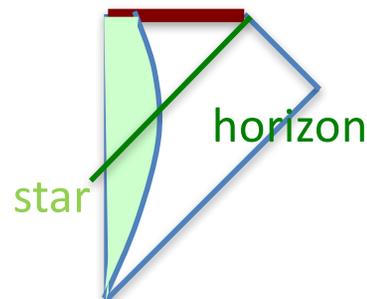
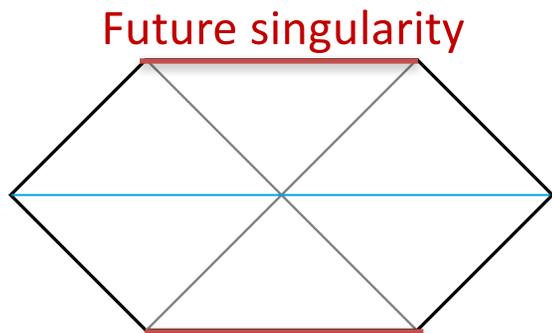
Eddington 1924 finds a non singular coordinate system but did not recognize (or comment on its significance).

Lemaitre 1933 First published statement that the horizon is not singular.

Einstein Rosen 1935 (Still call it a ``singularity'')

Szekeres, Kruskal 59-60 coordinates cover the full spacetime

Wheeler Fulling 62 It is a wormhole !

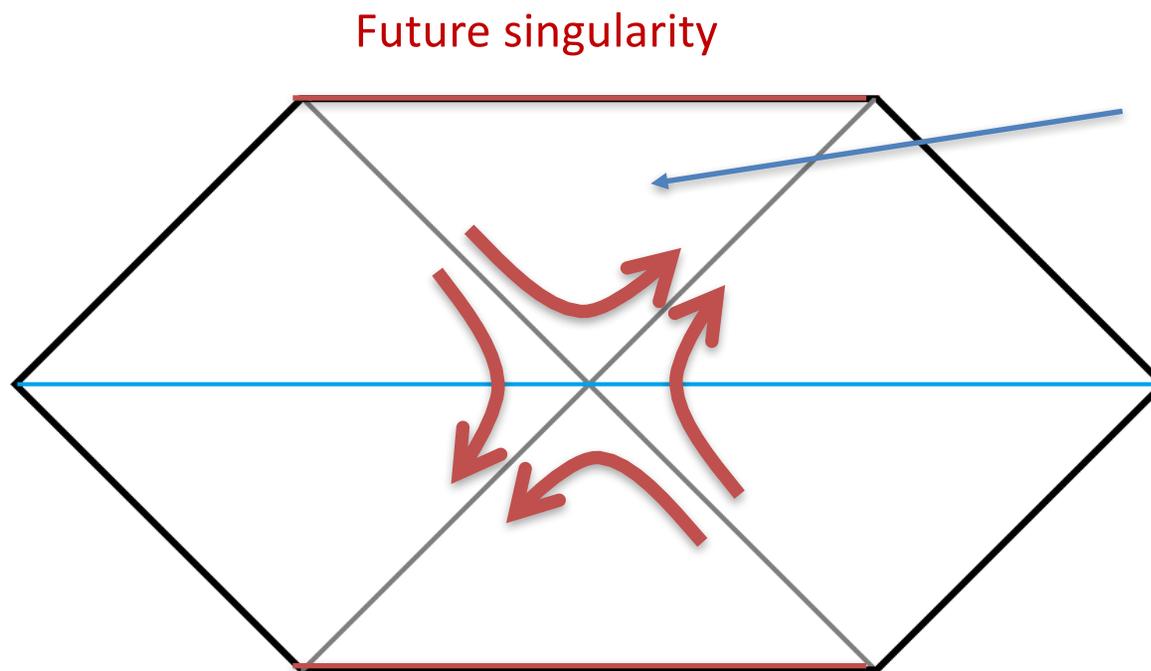


Oppenheimer  
Snyder 1939

Observer on the surface  
does not feel anything special  
at the horizon. Finite time to  
the singularity

It took 45 years to understand a classical solution...  
Why ?

The symmetries are realized in a funny way.  
The time translation symmetry  $\rightarrow$  boost symmetry at the  
bifurcation surface



Interior is time  
dependent for an  
observer falling in.  
It looks like a big crunch.

# Black holes in astrophysics

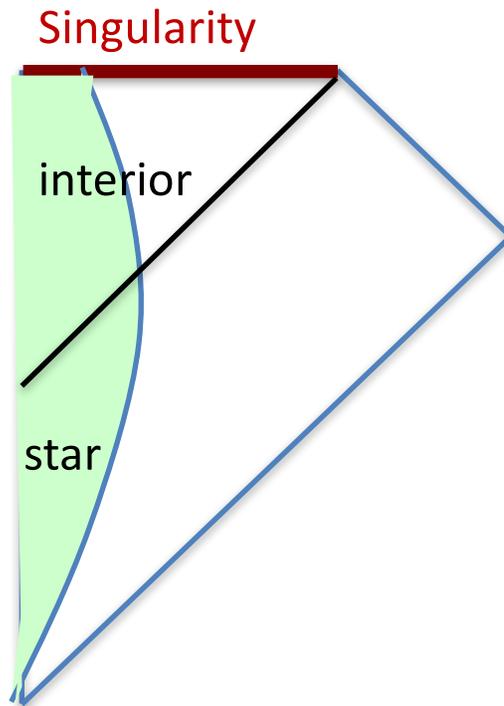
- Quasars (most efficient energy sources)
- Stellar mass black holes
- Sources for gravity waves !
- Our own supermassive black hole in the Milky Way...



# Interior and singularity

Quantum gravity is necessary at the singularity!

What signals ?



Singularity is behind the horizon.

It is shielded behind the black hole horizon that acts as a Schwarzschild.

# Black holes and quantum mechanics

- Black holes are one of the most surprising predictions of general relativity.
- Incorporating quantum mechanics leads to a new surprise:

Black holes are not black !

# Hot black holes

- Black holes have a temperature

$$T \sim \frac{\hbar}{r_H}$$

Hawking



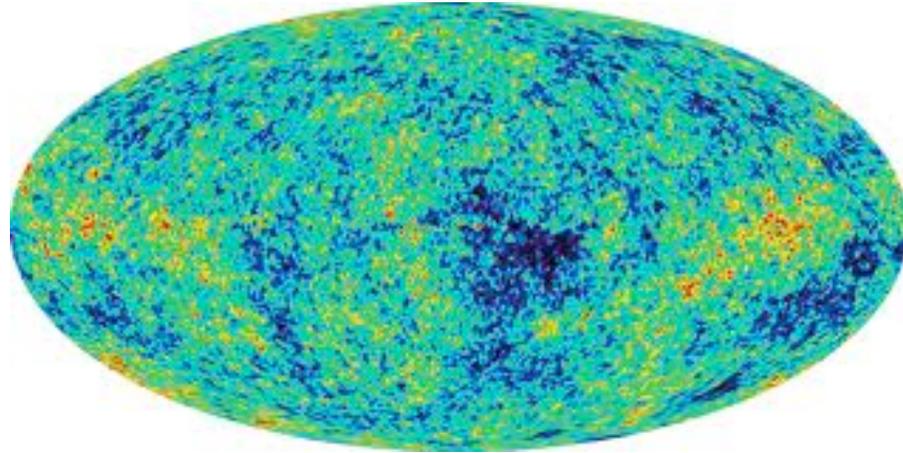
We can have white ``black holes''

- An accelerated expanding universe also has a temperature

$$T \sim \hbar H = \frac{\hbar}{R_H}$$

Chernikov, Tagirov,  
Figari, Hoegh-Krohn, Nappi,  
Gibbons, Hawking,  
Bunch, Davies, ....

Very relevant for us!



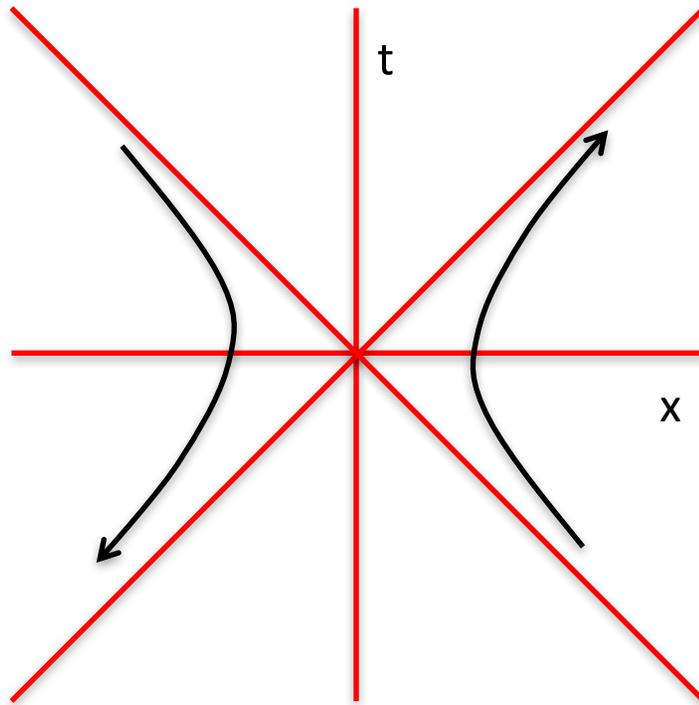
Quantum mechanics is crucial for understanding the large scale geometry of the universe.

# Why a temperature ?

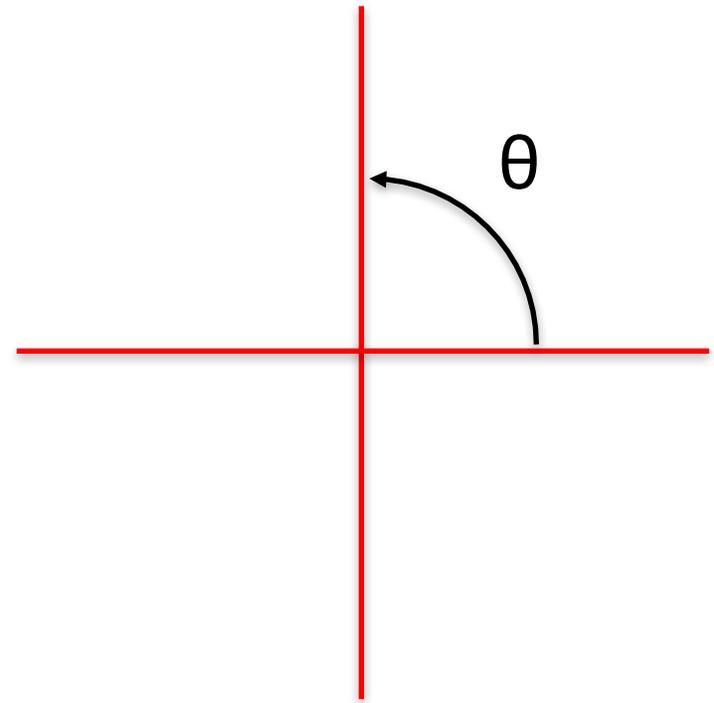
- Consequence of special relativity + quantum mechanics.

Flat space first

# Why a temperature ?



Lorentzian



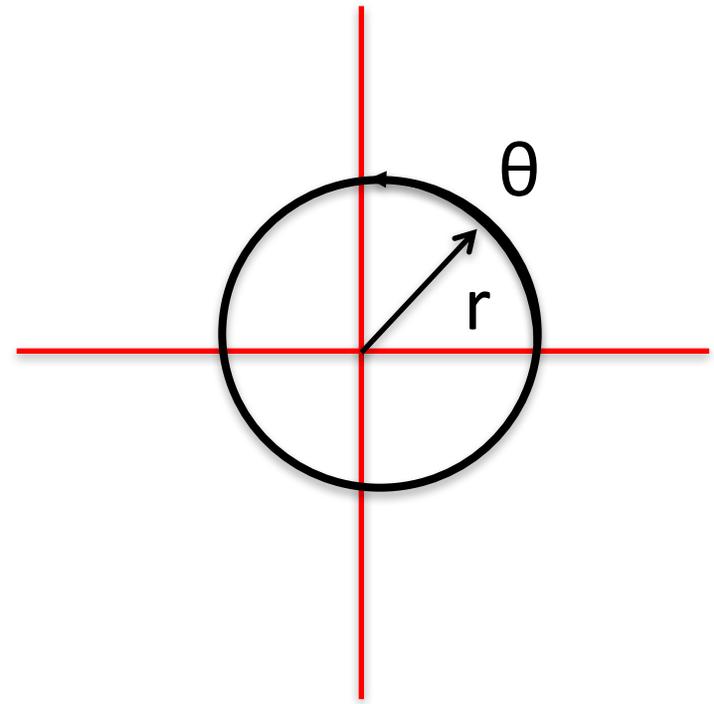
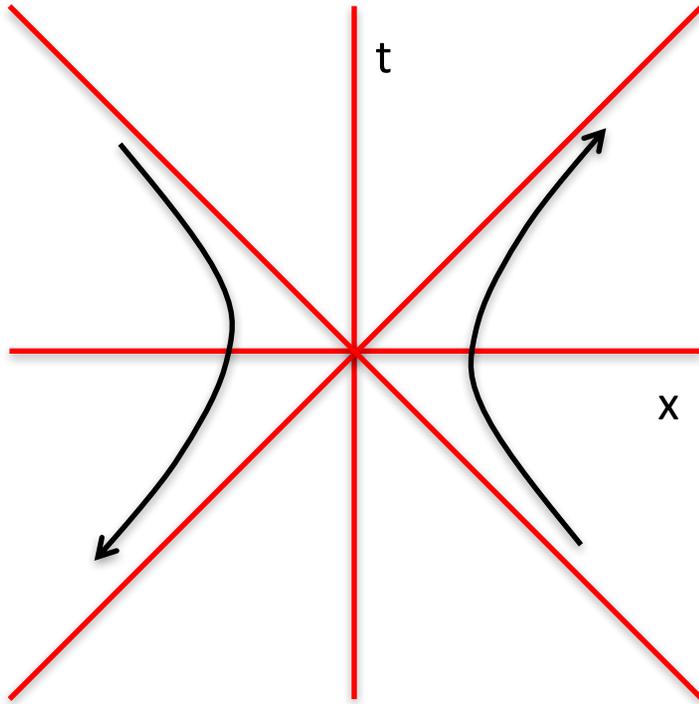
Euclidean

Accelerated observer  $\rightarrow$  energy = boost generator.

Time translation  $\rightarrow$  boost transformation

Continue to Euclidean space  $\rightarrow$  boost becomes rotation.

# Why a temperature ?



Continue to Euclidean space  $\rightarrow$  boost becomes rotation.

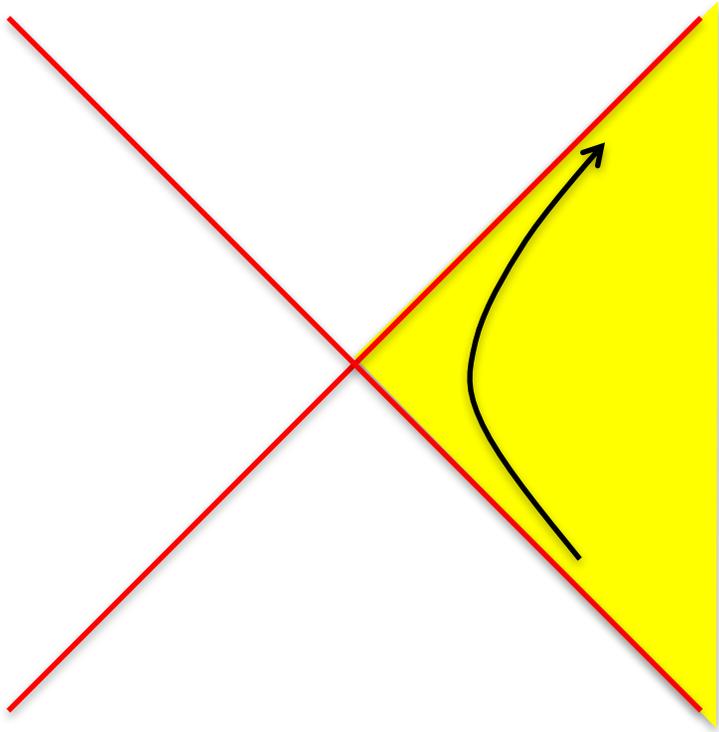
Angle is periodic  $\rightarrow$  temperature  $\beta = \frac{1}{T} = 2\pi r = \frac{2\pi}{a}$

Ordinary accelerations are very small,  $g = 9.8 \text{ m/s}^2 \rightarrow \beta = 1 \text{ light year}$

Thermische Unruhe

Bisognano Weichman,  $\uparrow$  Unruh

# Entanglement & temperature

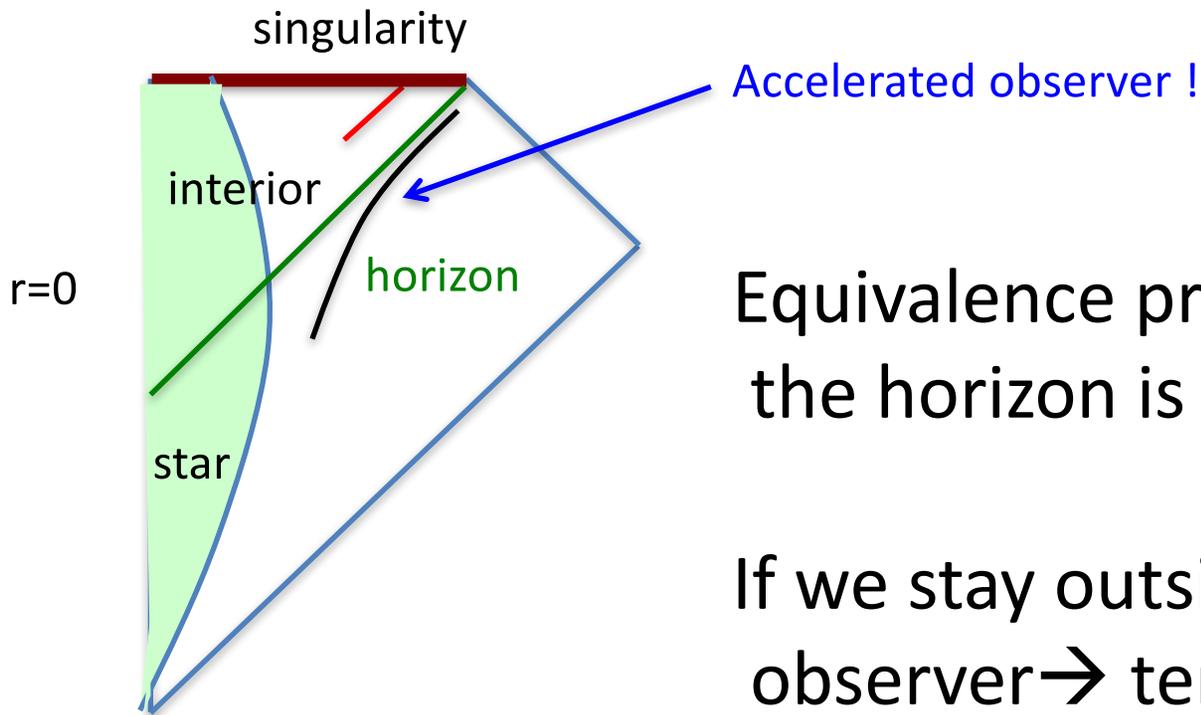


Horizon: accelerated observer only has access to the right wedge.

If we only make observations on the right wedge  $\rightarrow$  do not see the whole system  $\rightarrow$  get a mixed state (finite temperature).

General prediction, only special relativity + quantum mechanics + locality  
Vacuum is highly entangled !

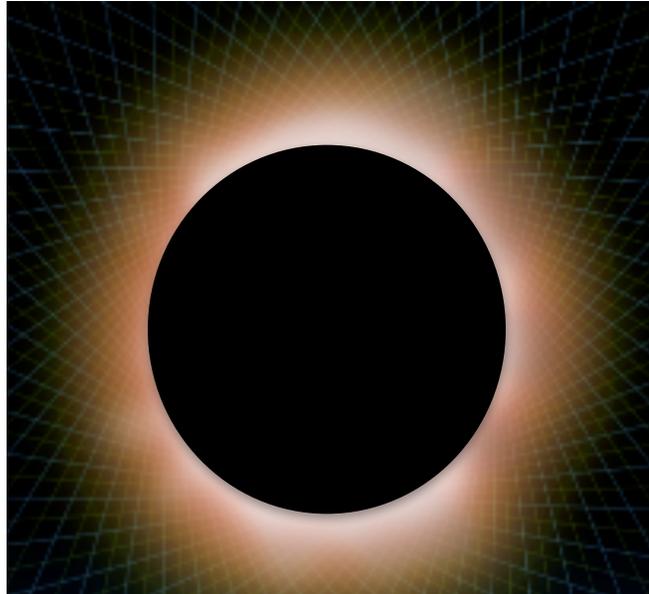
# Black hole case



Equivalence principle: region near the horizon is similar to flat space.

If we stay outside  $\rightarrow$  accelerated observer  $\rightarrow$  temperature.

Black hole from collapse



Black holes have a temperature

Do they obey the laws of thermodynamics ?

# Black hole entropy

$$T \sim \frac{\hbar}{r_H}$$

Special relativity near the black hole horizon

$$r_H \leftrightarrow M$$

Einstein equations

$$dM = TdS$$

1<sup>st</sup> Law of thermodynamics

$$S = \frac{(\text{Area})}{4\hbar G_N}$$

Black hole entropy

Bekenstein, Hawking

2<sup>nd</sup> Law  $\rightarrow$  area increase from Einstein equations and positive null energy condition. Hawking

# Including the quantum effects

$$S_{BH} = \frac{(\text{Area})}{4G_N} + S_{\text{entanglement}}$$



Entanglement entropy of quantum fields across the black hole horizon

Has been understood better in quantum field theory  
2<sup>nd</sup> Law extended to include this term.

Wall 2011

Bekenstein bound → automatic in relativistic quantum field theory.

Casini 2008

Focusing theorems and better understanding of the positivity of energy,  
and new "area" increase statements.

Bousso's talk  
Bousso, Englehardt, Wall,  
Faulkner, ...

# Bekenstein – Casini bound

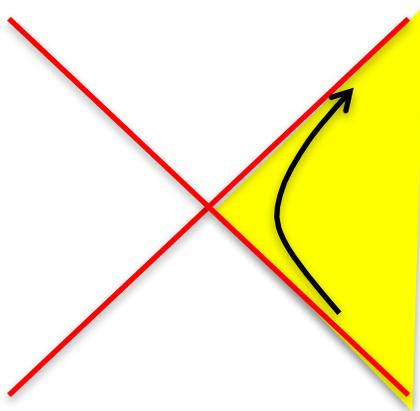
$$S \leq 2\pi ER$$

Bekenstein

When is this true ? Is it true ? Does it impose a constraint on QFT ?

$$\Delta S \leq 2\pi \Delta E_{\text{Rindler}}$$

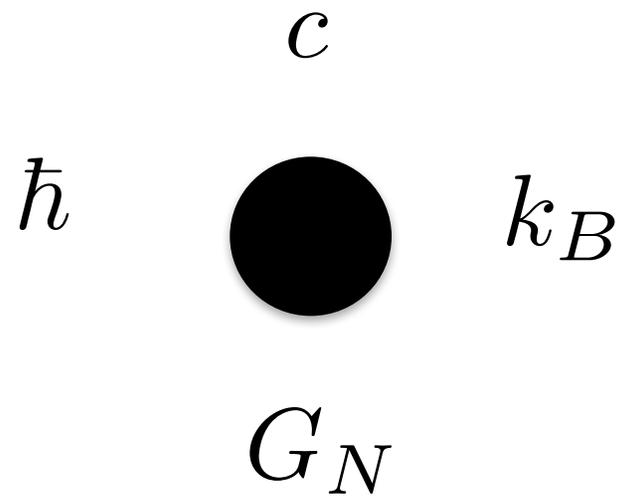
Casini 2008



It is always true in any relativistic QFT.

2<sup>nd</sup> Law always satisfied.

# General relativity and thermodynamics



# General relativity and thermodynamics

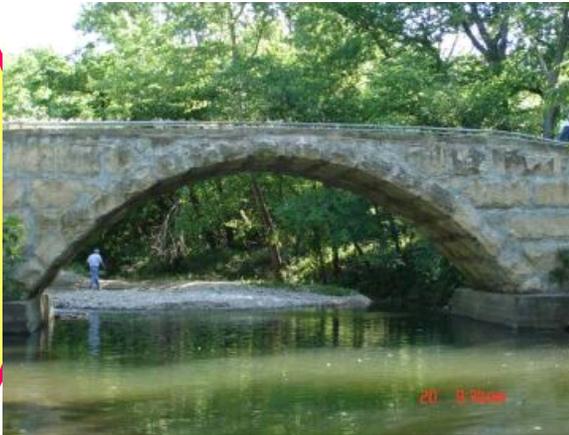
- Black hole seen from the outside = thermal system with finite entropy.
- Is there an exact description where information is preserved ?
- Yes...

# Gauge/Gravity Duality

(or gauge/string duality, AdS/CFT, holography)

Theories of quantum  
interacting particles

(very strongly interacting)



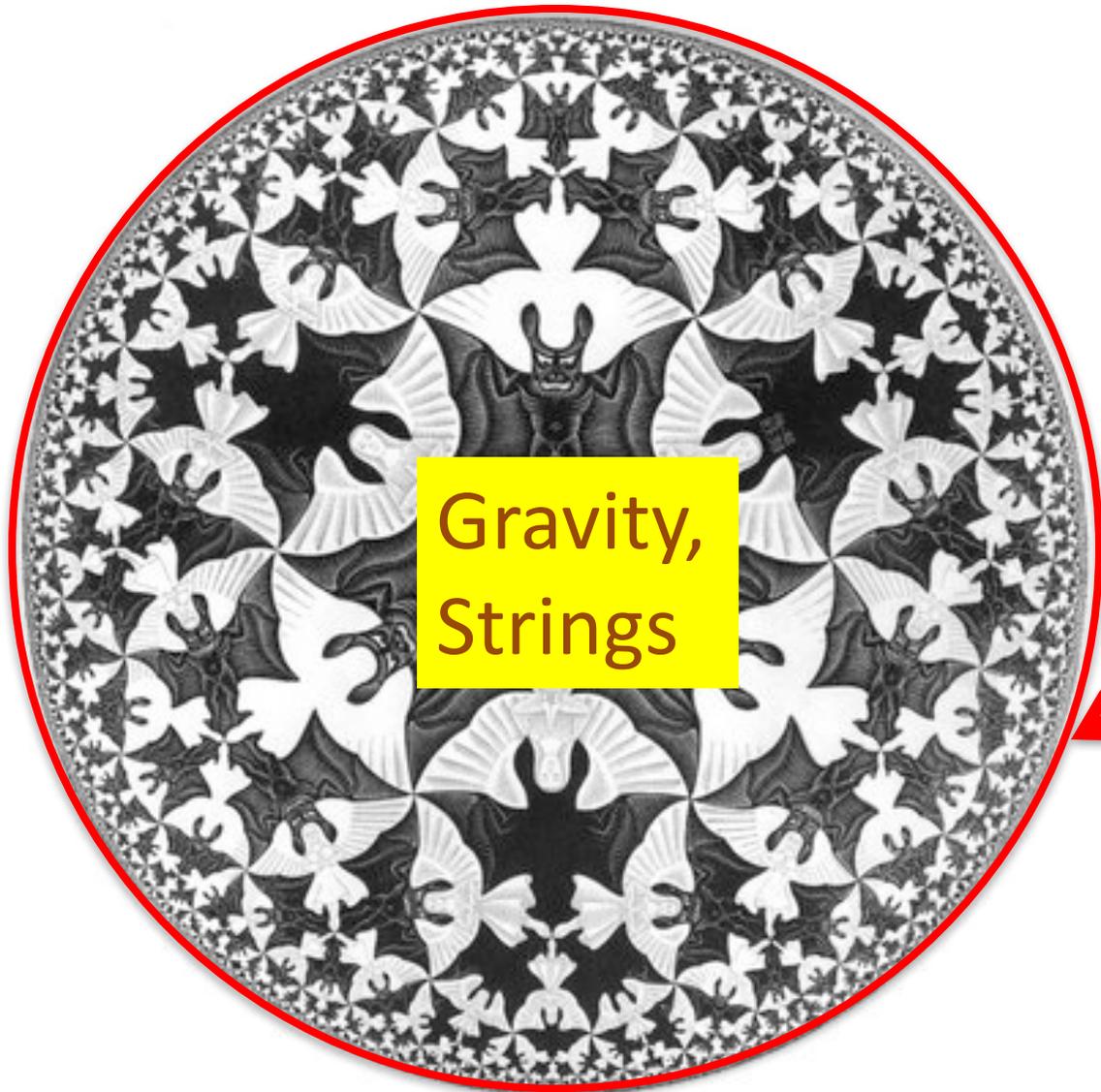
Quantum dynamical  
Space-time  
(General relativity)  
string theory

JM 97

Witten, Gubser, Polyakov, Klebanov

....

# Gravity in asymptotically Anti de Sitter Space



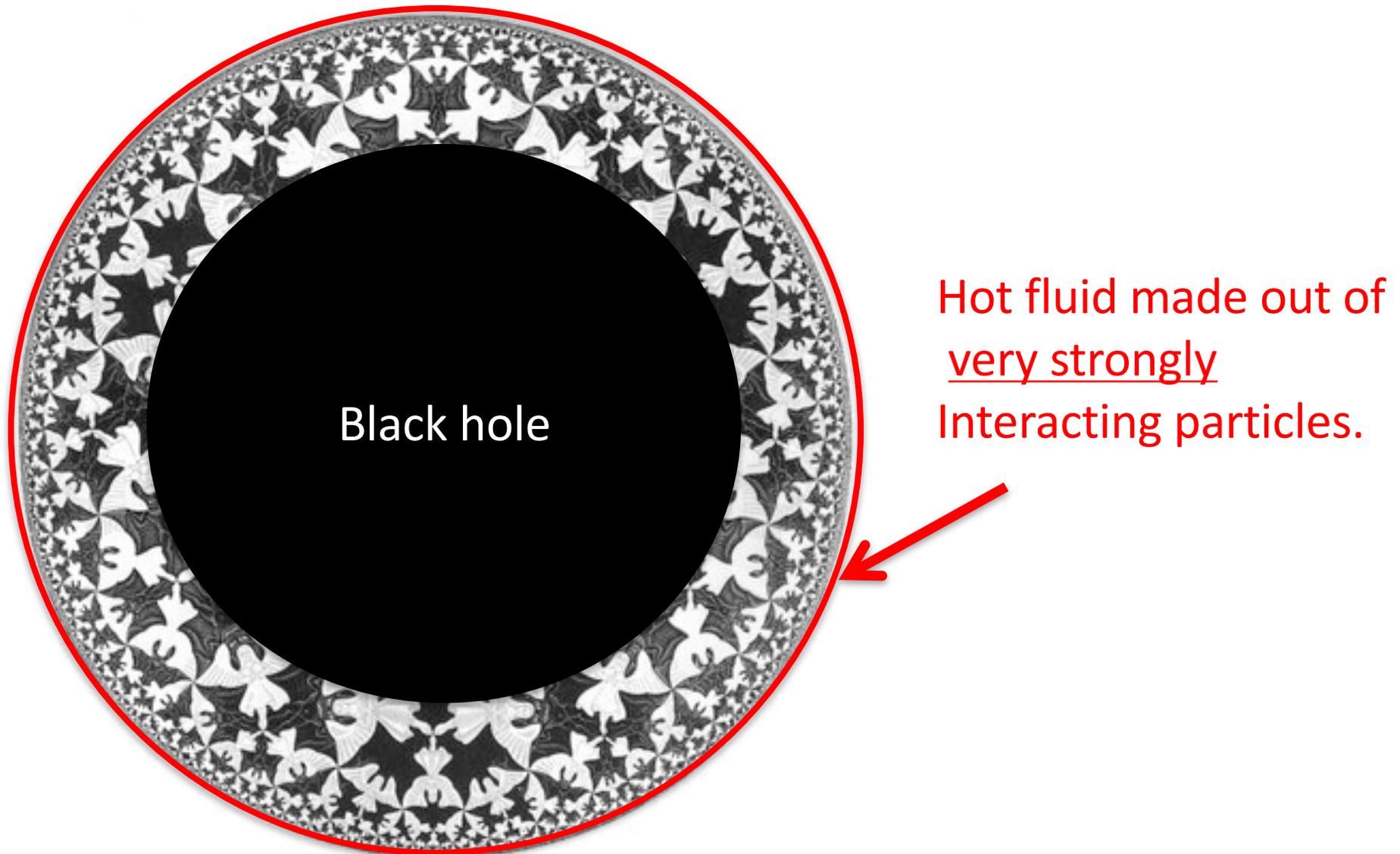
Gravity,  
Strings

Duality

Quantum interacting particles  
quantum field theory



# Black holes in a gravity box



# Lessons for black holes

- Black holes as seen from outside (from infinity) are like an ordinary quantum system.
- Black hole entropy = ordinary entropy of the quantum system.
- Absorption into the black hole = thermalization
- Chaos  $\rightarrow$  near horizon gravity

Shenker Stanford  
Kitaev , 2013-2015

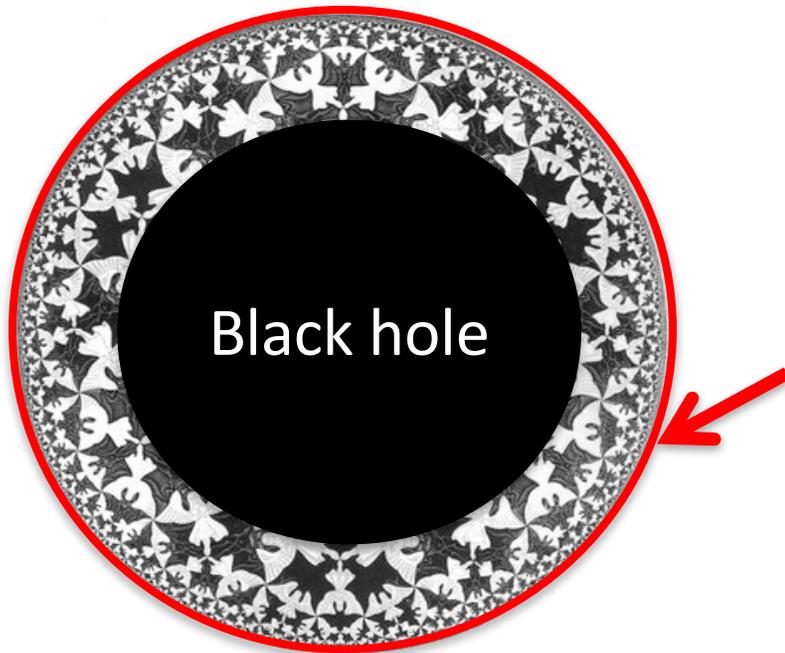
- Interior ?

Mathur, Almheiri, Marolf, Polchinski, Sully

In the meantime

# Black holes as a source of information

## Black holes as toy models



Used to model strongly interacting systems in high energy physics or condensed matter physics.

Hot fluid made out of strongly  
Interacting particles.

Key insights into the theory of hydrodynamics  
with anomalies.

Damour, Herzog, Son, Kovtun, Starinets, Bhattacharyya, Hubeny, Loganayagam,  
Mandal, Minwalla, Morita, Rangamani, Reall, Bredberg, Keeler, Lysov, Strominger...

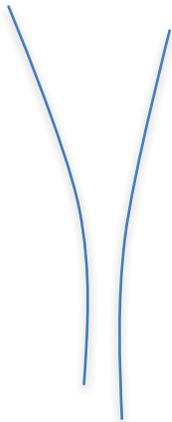
Let us go back to chaos



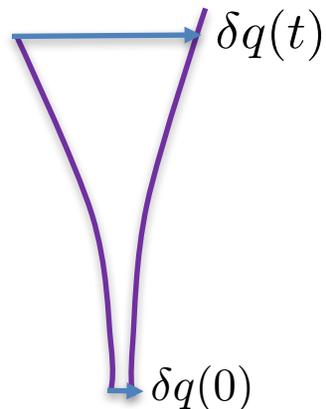
Chaos → divergence of nearby trajectories

Thermal system → average over all trajectories

Growth → Where you are after the perturbation  
vs. where you would have been.



Classical



$$\delta q(t) \propto e^{\lambda t}$$

$$\frac{\delta q(t)}{\delta q(0)} = \{q(t), p(0)\} , \quad \{q(t), p(0)\}^2$$

Quantum

$$[Q(t), P(0)]^2$$

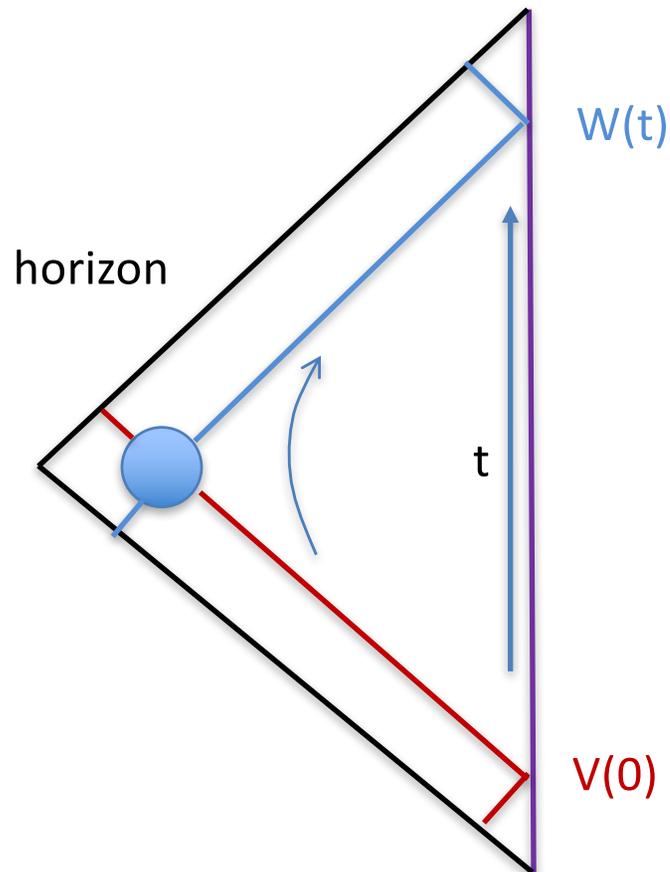
Quantum General:

$$\langle [W(t), V(0)]^2 \rangle_{\beta} \propto \frac{1}{N} e^{\lambda t}$$

$W, V$  are two "simple" (initially commuting) observables.

Imagine we have a large  $N$  system. This is the definition of the quantum Liapunov exponent

For quantum systems that have a gravity dual



Commutator  $\rightarrow$  involves the scattering amplitude between these two excitations.

Leading order  $\rightarrow$  graviton exchange

Large  $t \rightarrow$  large boost between the two particles.

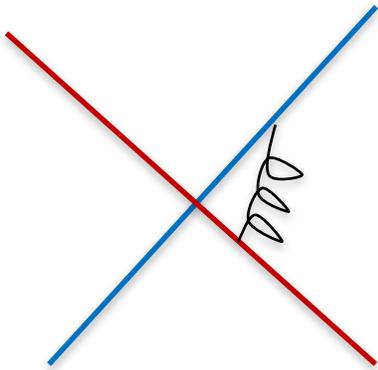
Gravitational interaction has spin 2,  
Shapiro time delay proportional to energy.

Energy goes as  $e^t$

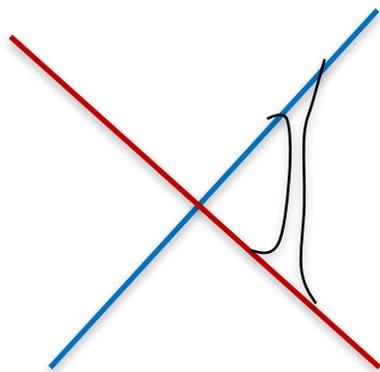
$$\langle [W(t), V(0)]^2 \rangle_\beta \propto \frac{1}{N} e^{\frac{2\pi}{\beta} t}$$

$$\lambda = \frac{2\pi}{\beta} = 2\pi T$$

# Can it be different ?



Graviton  $\rightarrow$  phase shift :  $\delta(s) \sim G_N s \longrightarrow \lambda = \frac{2\pi}{\beta}$



String  $\rightarrow$  phase shift

$$\delta(s) \sim G_N s^{1+\alpha' t} \longrightarrow \lambda = \frac{2\pi}{\beta} \left(1 - \frac{l_s^2}{R^2}\right)$$

Typical size of string  
(of graviton in string theory)



$s, t =$  Mandelstam invariants

Radius of curvature of black hole

It can be less...

More ?

In flat space a phase shift has to scale with a power of  $s$  less than one in order to have a causal theory

Maybe there is a bound...

# Black holes as the most chaotic systems

$$\lambda \leq \frac{2\pi}{\beta} = 2\pi T$$

Sekino Susskind

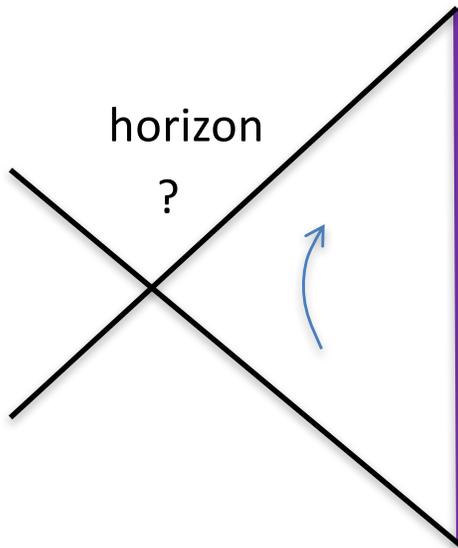
JM, Shenker, Stanford

For any large N (small  $\hbar$ ) quantum system.

(Strings connect weakly coupled to strongly coupled systems)

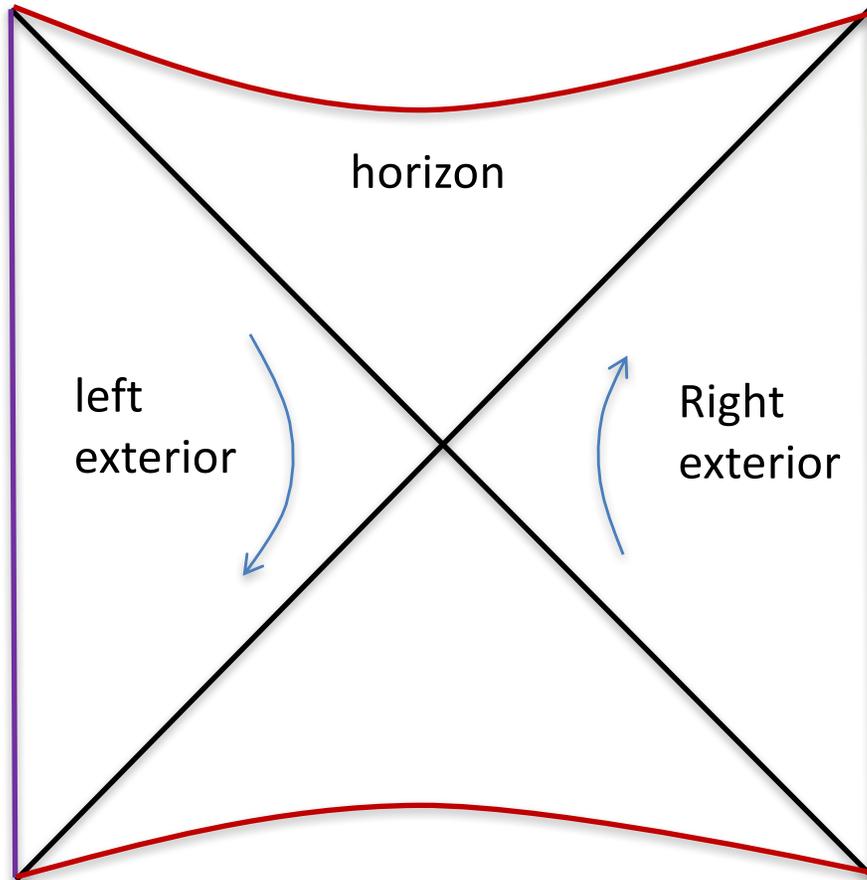
# How do we get order from chaos ?

## How do we get the vacuum from chaos, or from a chaotic quantum system?



Example: hydrodynamics, we get something simple for some interactions, but it is more complicated with very small interactions (Boltzman equation).

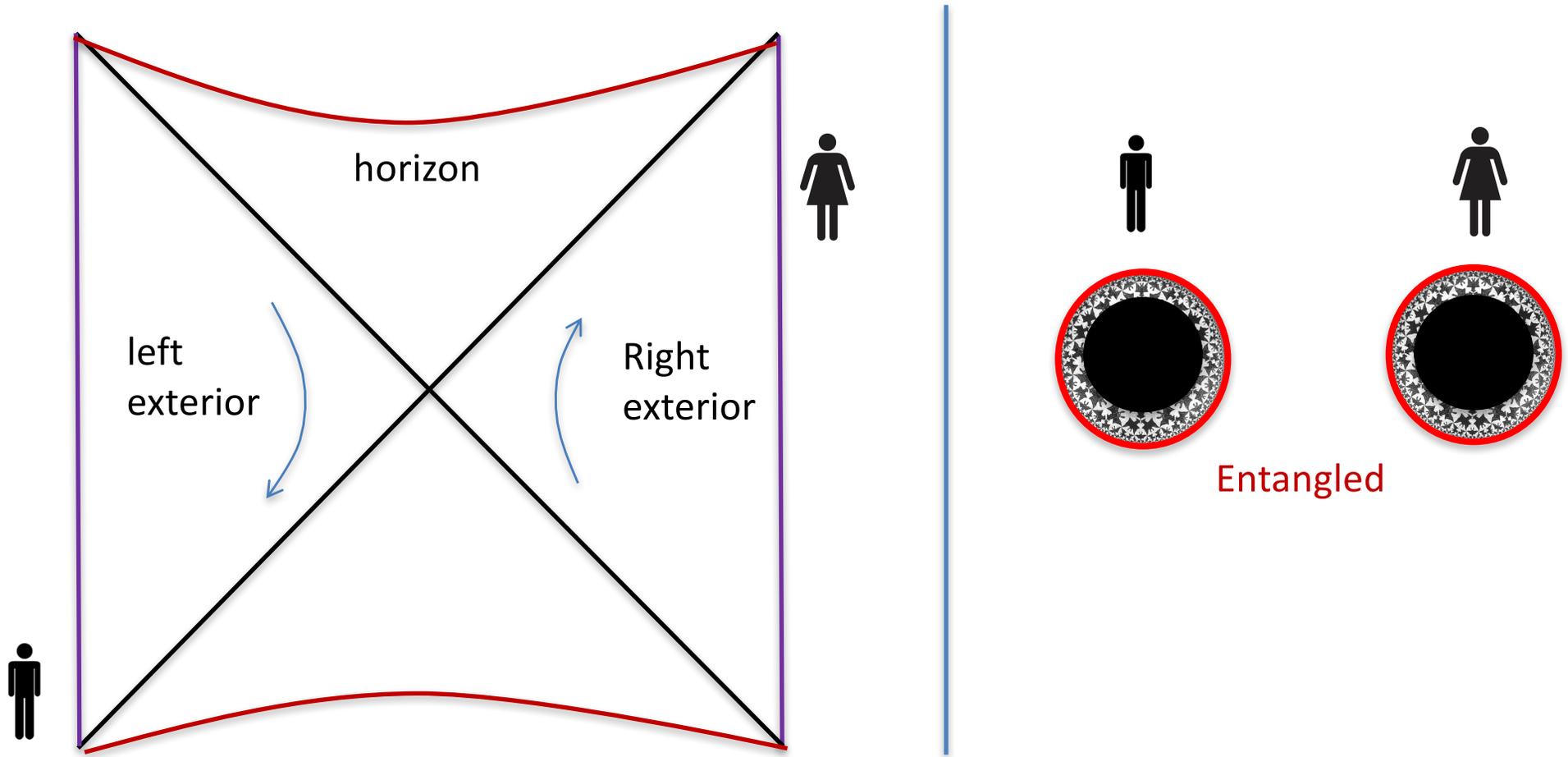
# The full Schwarzschild wormhole



No need to postulate any exotic matter

No matter at all !

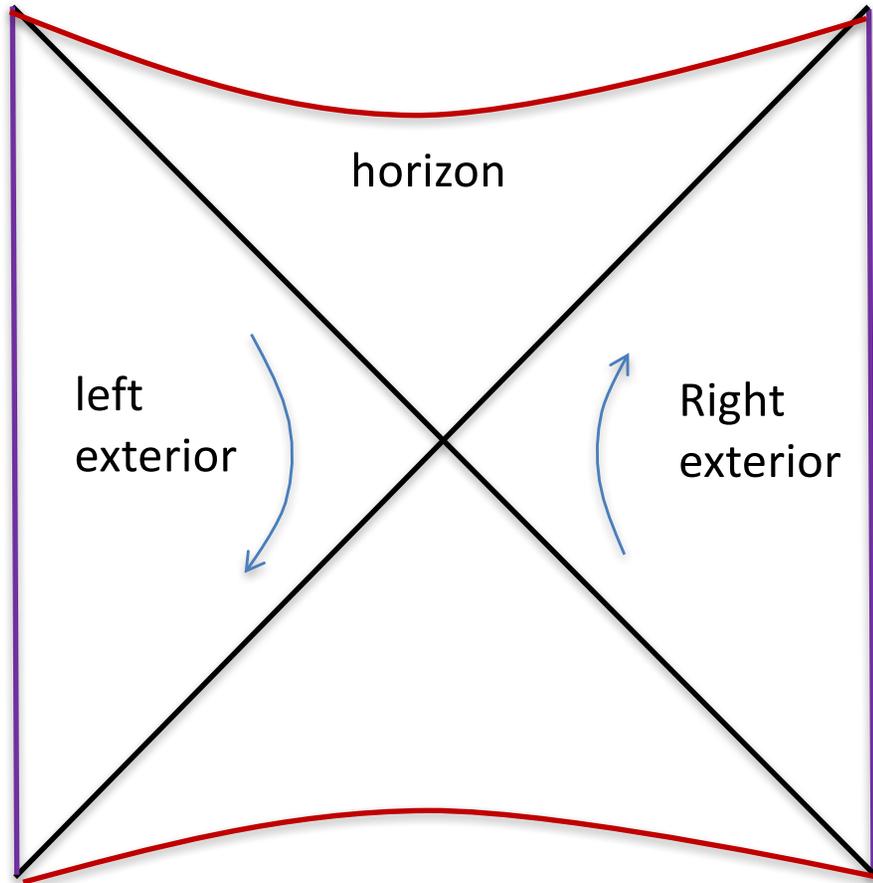
# View it as an entangled state



$$|TFD\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L |E_n\rangle_R$$

Israel 70's  
JM 00's

# Symmetry



What is this funny  
“time translation symmetry” ?

$$U = e^{it(H_R - H_L)}$$

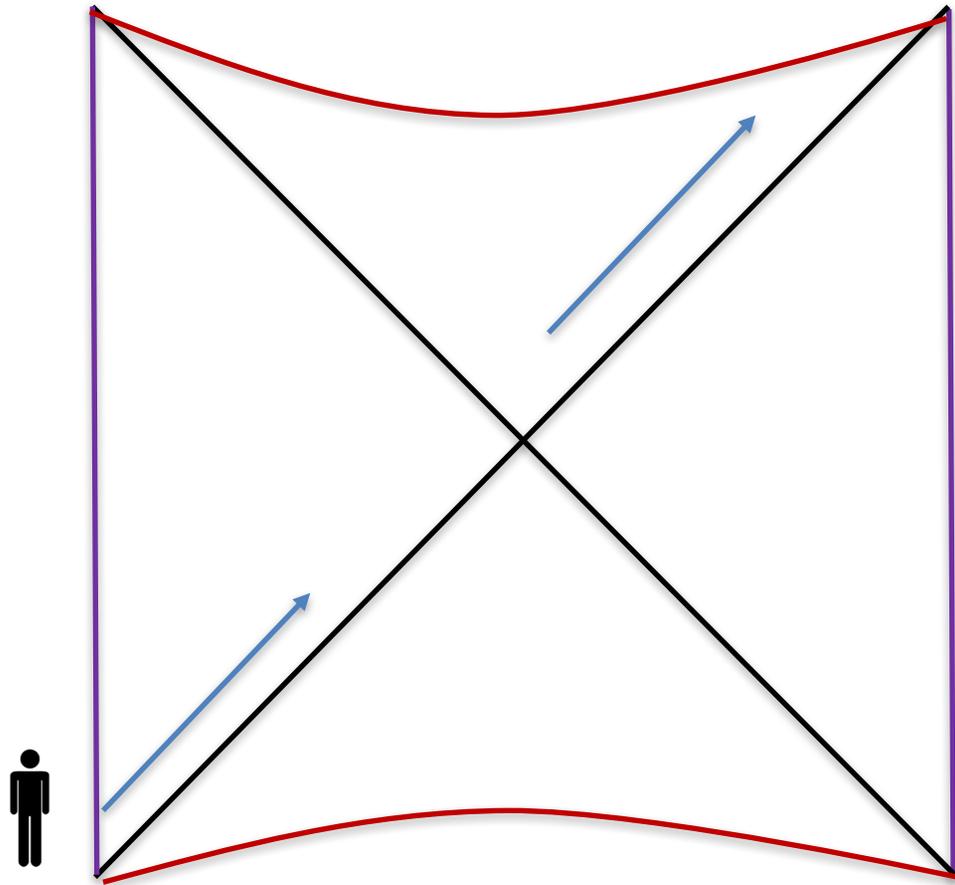
Exact symmetry

Exact boost symmetry!

This symmetry should be exact  
for any theory of quantum gravity

$$|TFD\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L |E_n\rangle_R$$

# True causal separation



If Bob sends a signal , then Alice cannot receive it.

These wormholes are not traversable

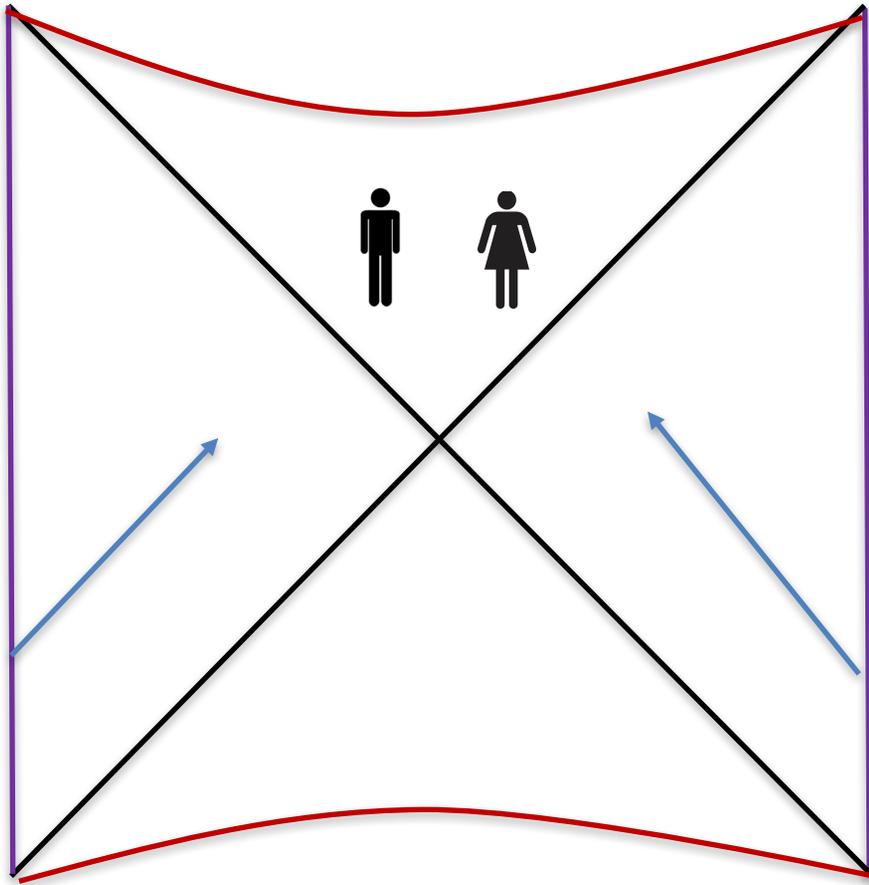
Relay on Integrated null energy condition  
(Now a theorem, proven using entanglement methods) [Balakrishnan, Faulkner, Khandker, Wang](#)

Not good for science fiction.

Good for science!

$$|TFD\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L |E_n\rangle_R$$

# Interior is common

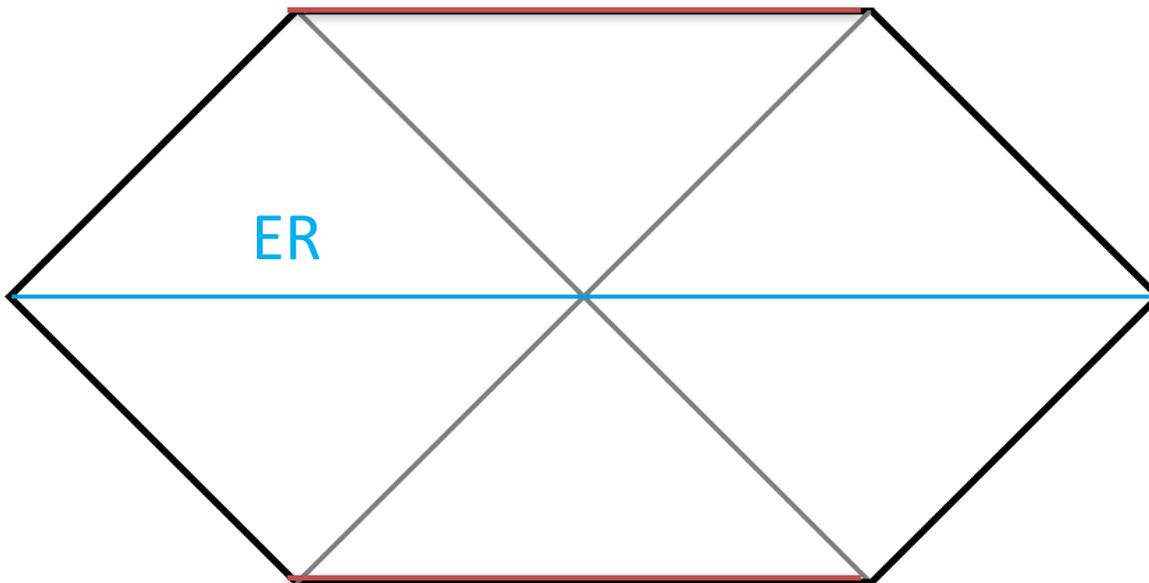


If they jump in,  
they can meet in the interior !

But they cannot tell anyone.

$$|TFD\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L |E_n\rangle_R$$

# Spacetime connectivity from entanglement



$$ER = EPR$$

Van Raamsdonk  
Verlinde<sup>2</sup>  
Papadodimas Raju

JM Susskind

# Entanglement and geometry



Local boundary quantum bits are highly interacting and very entangled

Ryu, Takayanagi, Hubeny, Rangamani

$$S(R) = \frac{A_{\min}}{4G_N}$$

Generalization of the Black hole entropy formula.

Interesting connections to quantum information theory

Quantum error correction

Almheiri, Dong, Harlow, Preskill, Yoshida, Pastawski

Complexity theory

Harlow, Hayden, Brown, Susskind

# Conclusions

- Black holes are extreme objects: most compact, most efficient energy conversion, most entropic, most chaotic,...
- Most confusing...
- The process of unravelling these confusions has lead to better understand of gravity, quantum systems, string theory and their interconnections.
- Black holes are not only in the cosmos, but can also be present in the lab.
- And there are still very important confusions and open problems: Interior and singularity ?

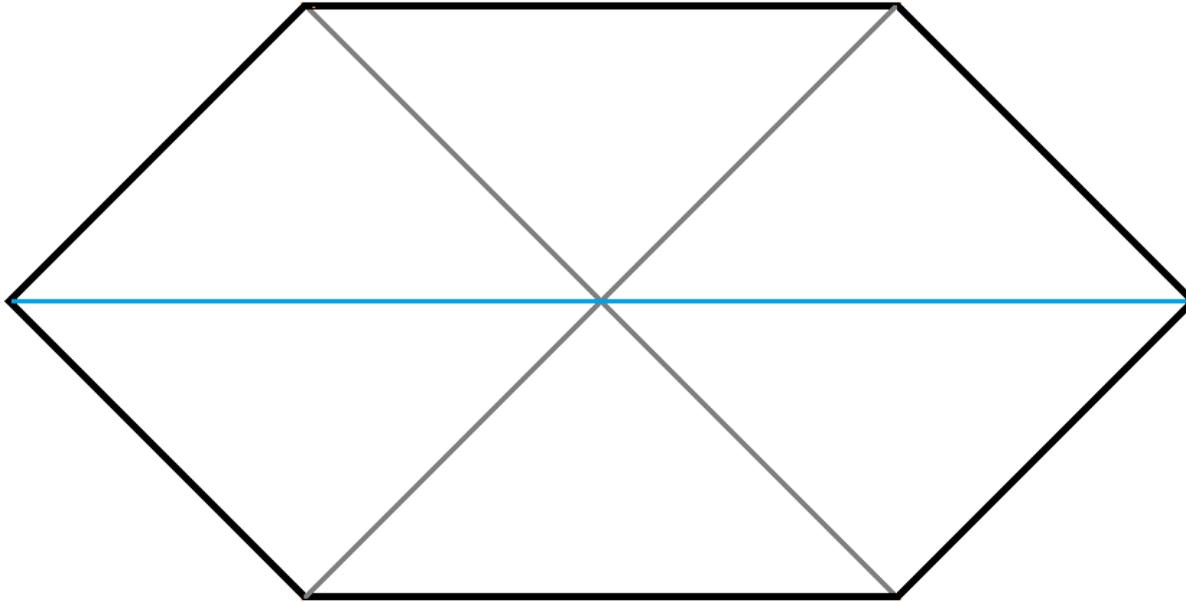
Thank you!





Extra slides

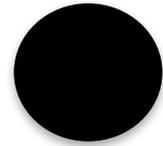
# Full Schwarzschild solution



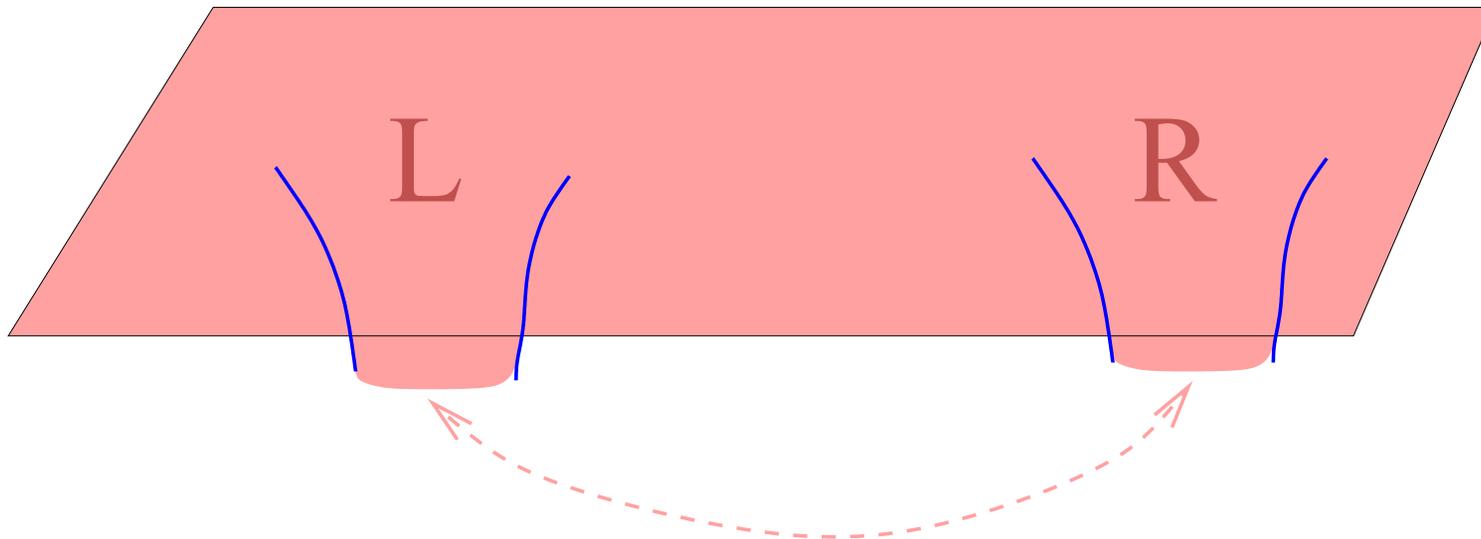
Eddington, Lemaitre,  
Einstein, Rosen,  
Finkelstein  
Kruskal

Simplest spherically symmetric solution of pure Einstein gravity  
(with no matter)

# Wormhole interpretation.

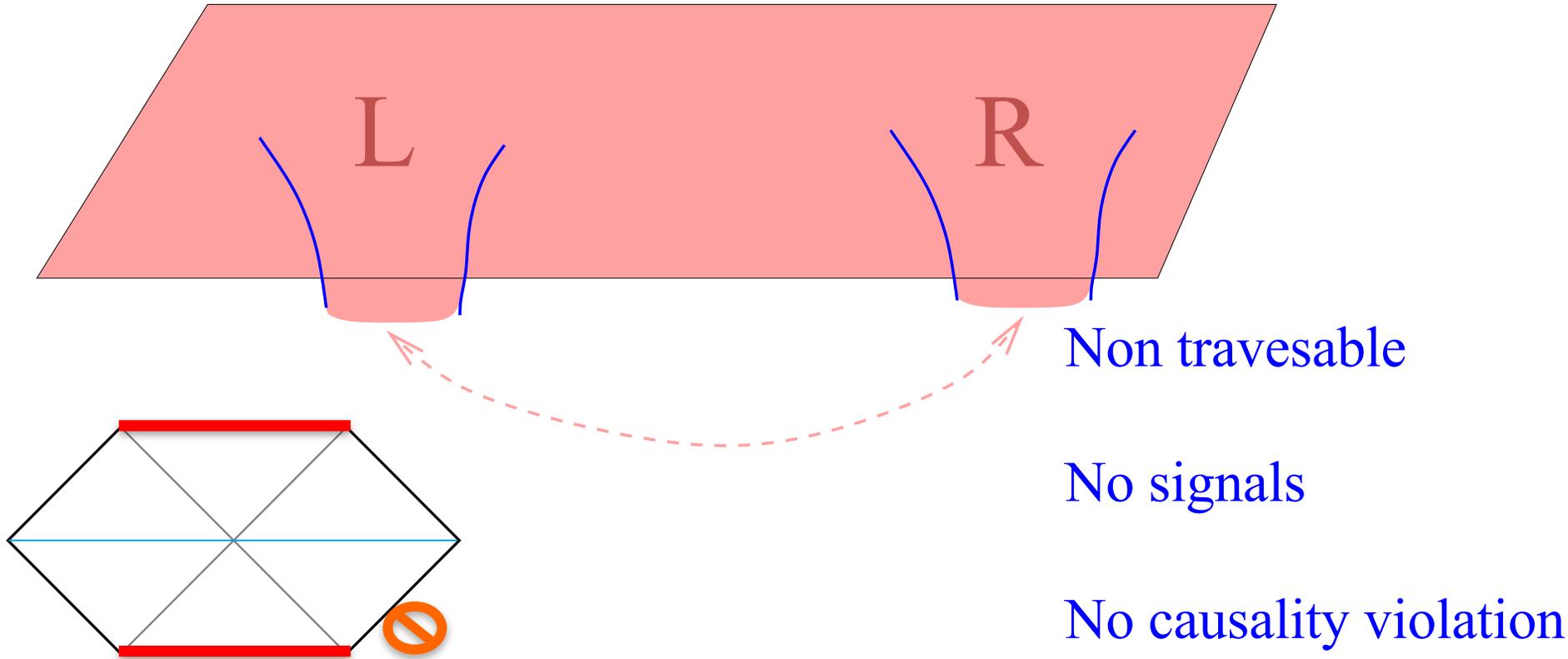


# Wormhole interpretation.

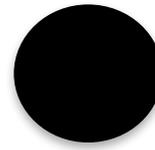


Note: If you find two black holes in nature, produced by gravitational collapse, they will not be described by this geometry

# No faster than light travel



Fuller, Wheeler, Friedman, Schleich, Witt, Galloway, Wooglar



In the exact theory,  
each black hole is described by a set of microstates from the outside

Wormhole is an entangled state

$$|\Psi\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L \times |E_n\rangle_R$$

EPR

Israel  
JM

Geometric connection  
from entanglement. ER = EPR

Susskind JM

Stanford, Shenker, Roberts, Susskind

MAY 15, 1935

PHYSICAL REVIEW

VOLUME 47

## Can Quantum-Mechanical Description of Physical Reality Be Considered Complete?

A. EINSTEIN, B. PODOLSKY AND N. ROSEN, *Institute for Advanced Study, Princeton, New Jersey*

(Received March 25, 1935)

EPR

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JULY 1, 1935

PHYSICAL REVIEW

VOLUME 48

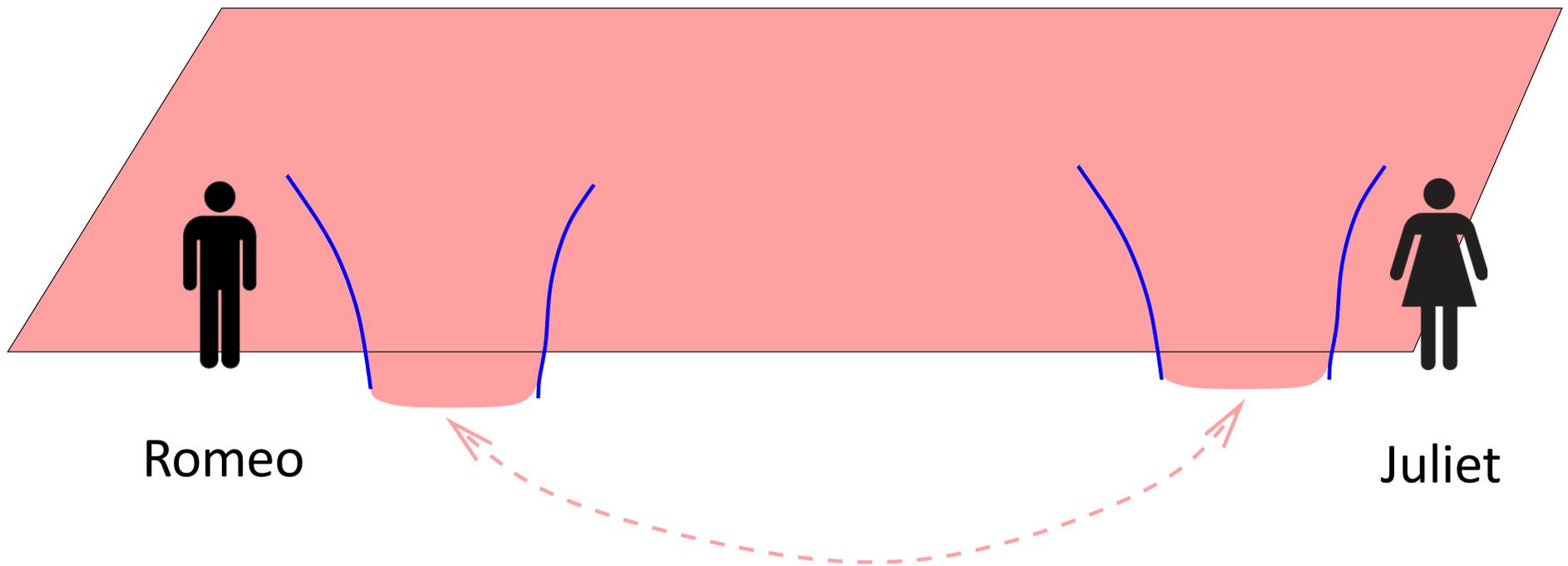
## The Particle Problem in the General Theory of Relativity

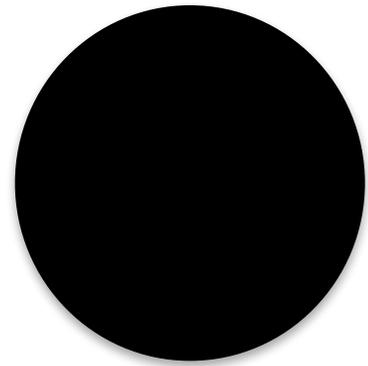
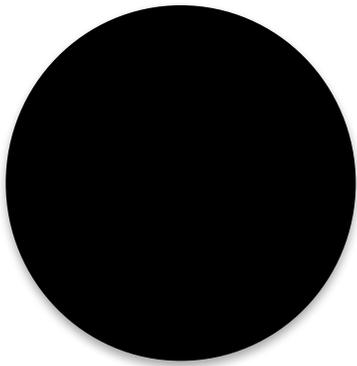
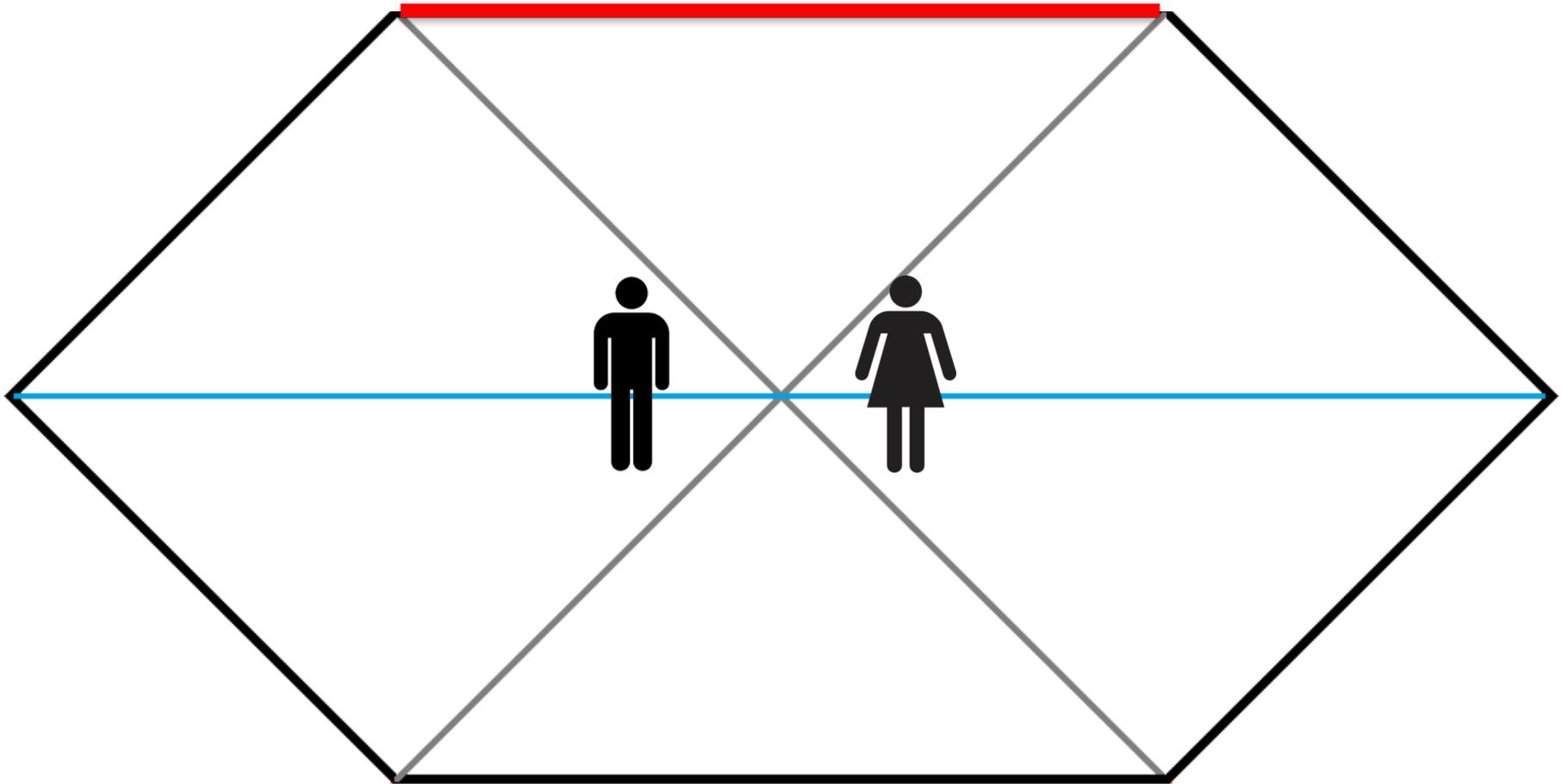
A. EINSTEIN AND N. ROSEN, *Institute for Advanced Study, Princeton*

(Received May 8, 1935)

ER

# A forbidden meeting





# String theory

- String theory started out defined as a perturbative expansion.
- String theory contains interesting solitons: D-branes.
- Using D-branes one can “count” the number of states of extremal charged black holes in certain superstring theories. Polchinski
- D-branes inspired some non-perturbative definitions of the theory in some cases. Strominger Vafa

Matrix theory: Banks, Fischler, Shenker, Susskind

Gauge/gravity duality: JM, Gubser, Klebanov, Polyakov, Witten

# Entanglement and geometry

- The entanglement pattern present in the state of the boundary theory can translate into geometrical features of the interior.  
Van Raamsdonk
- Spacetime is closely connected to the entanglement properties of the fundamental degrees of freedom.
- Slogan: Entanglement is the glue that holds spacetime together...
- Spacetime is the hydrodynamics of entanglement.



# Questions

- Black holes look like ordinary thermal systems if we look at them from the outside. We even have some conjectured exact descriptions.
- How do we describe the interior within the same framework that we describe the exterior ?  
Modern version of the information paradox;  
Mathur, Almheiri, Marolf, Polchinski, Stanford, Sully,..
- Once we figure it out: what is the singularity ?
- What lessons do we learn for cosmology ?